Effects of baryon number density fluctuation around QCD critical point Kazuhiko Kamikado (RIKEN)

Based on Prog. Theor. Exp. Phys. (2013) 053D01, K. Kamikado, T. Kunihiro, K. Morita, and A. Ohnishi

Motivation

QCD phase diagram



Critical region

contour of quark number susceptibility



Landau potential at finite density

$$\begin{split} U &= a\sigma^2 + b\sigma^4 + c\sigma^6 - d\sigma \quad (\mu = 0) \\ & \sigma \sim \bar{\psi}\psi \text{(chiral order parameter)} \end{split}$$



1st order boundary and critical point (2nd order)

diverge near the critical point. $\langle (\bar{\psi}\psi)^2 \rangle, \ \langle (\bar{\psi}\gamma_0\psi)^2 \rangle \gg 1$

We need to evaluate the size or form of the critical region.

At finite μ , coupling between σ and quark-density (phonon) mode is inevitable.

Formalism

Quark-meson model (ψ , σ and π) with density fluctuation (ϕ)

 g_d determines the strength of the coupling between σ and ϕ .

Functional renormalization group equation

$$k\partial_{k}\Gamma_{k}[\varphi] = \frac{1}{2}\operatorname{Tr}\left[\frac{k\partial_{k}R_{kB}}{R_{kB} + \Gamma_{k}^{(0,2)}[\varphi]}\right] - \operatorname{Tr}\left[\frac{k\partial_{k}R_{kF}}{R_{kF} + \Gamma_{k}^{(2,0)}[\varphi]}\right] \qquad \mathsf{Sc}$$
C. Wetterich (1993)
R is arbitrary cutoff function. Our choices are

$$R_{kB} = (k^{2} - \vec{p}^{2})\theta[k^{2} - \vec{p}^{2}], \ R_{kF} = ik\frac{\vec{p}}{|\vec{p}|}\theta[k^{2} - \vec{p}^{2}] \qquad \Gamma$$
D. Litim (2000)

cale dependent effective action $r_{k=\Lambda}[\phi] = S[\phi]$ UV: classical $\Gamma_{k=0}[\phi] = \Gamma[\phi]$ IR: quantum

20 20 Local potential approximation (LPA) $\Gamma_{k}^{LPA} = \frac{Z_{\sigma k}}{2} \partial \sigma \partial \sigma + \frac{Z_{\pi k}}{2} \partial \vec{\pi} \partial \vec{\pi} + \frac{Z_{\varphi k}}{2} \partial \varphi \partial \varphi \partial \varphi_{\mu[MeV]}^{300} U_{k}(\sigma^{2} + \pi^{2}, \varphi) + c\sigma$ 260 $Z_{\sigma k} = Z_{\pi k} = Z_{\varphi k} = 1$ $g_{d} = 0.2g_{s}$ 80 We solve functional differential equation for the scale $\sim Z_{\varphi} p^2$ dependent potential U_k in $\frac{1}{2}he^{\omega}2D$ grid (σ_n, ϕ_m) space. MeV] flows of wave funtetion 40 renormalization are neglected 2020 $U_k(\sigma_n,\phi_m)$ 320 340 280 240 300 260 µ[MeV]

80

Conditions for the order parameters:

Results





Behaviors of the soft mode mass at $T = T_c$





M =

 $\frac{\partial^2 U}{\partial \sigma \partial \varphi} \\ \frac{\partial^2 U}{\partial^2 U}$

 $\overline{\partial arphi \partial arphi}$,

 σ_{0}, φ_{0}

 $\begin{bmatrix} \frac{\partial^2 U}{\partial \sigma \partial \sigma} \\ \frac{\partial^2 U}{\partial \sigma^2 U} \end{bmatrix}$

 $\overline{\partial \varphi \partial \sigma}$

 $\frac{\partial U_{k=0}}{\partial \sigma} = \frac{\partial U_{k=0}}{\partial \varphi} = 0$

Contour of smaller mass M





	210	200	200	500	J
μ [MeV]			µ[MeV]		

- A kind of level repulsion between σ and ϕ is founded.
- This level repulsion enhances the softening behavior of M₋.

Summary

- We have studied the QCD critical point and its critical region by using the functional-RG method.
- We have introduced a model which includes the baryon-density fluctuation mode as well as the chiral condensate.
- We have seen an expansion of the critical region with the coupling increasing.
- The expansion behavior is caused by a level repulsion of the masses. Then the inclusion of a new mode will generally cause the expansion of the critical region.