# Differences and similarities between fundamental and adjoint fermions in SU(N) gauge theories

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## Aim

This work aim at clarifying what is the most essential difference between fundamental and adjoint fermions.

This subject will be talked by Dr. Kashiwa.

This poster is based on Phys. Rev. D88 (2013) 016002 [arXiv:1304.3274].

## Our question and its answer

What is the most essential difference between fundamental and adjoint quarks?



The answer is  $Z_3$  symmetry. Adjoint quarks preserve  $Z_3$  symmetry, but fundamental quarks break the symmetry.



To understand this point clearly, we consider the new boundary condition for fundamental quarks that preserve  $Z_3$  symmetry.

# QCD action with fundamental quarks

$$S_0 = \int d^4x \left[ \sum_f \bar{q}_f (\gamma_\nu D_\nu + m_f) q_f + \frac{1}{4g^2} F_{\mu\nu}^{a^2} \right],$$

The action is invariant under Z<sub>3</sub> transformation

$$q \rightarrow Uq$$
,  $A \rightarrow UAU^{-1} - \frac{i}{g} (\partial U)U^{-1}$ ,  
 $U(x,\beta) = \exp[i2\pi/3]U(x,0)$ 

But the fermion boundary condition is not.

### Boundary condition of fundamental fermion

$$q_f(x,\beta=1/T) = -\exp[i\theta]q_f(x,0)$$
Twist angle
$$Z_3 \text{ transformation}$$

$$\theta = 0$$

$$\theta = 2\pi/3$$

 $Z_3$  transformation rotates the twist angle by  $2\pi/3$ .

#### Construction of fundamental quarks with Z<sub>3</sub> symmetry

#### Consider degenerate three-flavor system

$$N_c = N_f = 3, \quad m_1 = m_2 = m_3$$

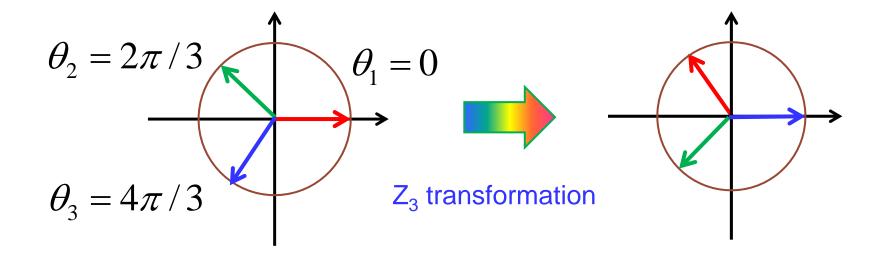
#### QCD with FTBC has $Z_3$ symmetry.

Flavor-dependent twist boundary condition (FTBC):

$$q_1(x, \beta = 1/T) = -q_1(x, 0),$$
  
 $q_2(x, \beta = 1/T) = -\exp[i2\pi/3]q_2(x, 0),$   
 $q_3(x, \beta = 1/T) = -\exp[i4\pi/3]q_3(x, 0).$ 

## Fundamental quarks with FTBC

$$q_f(x, \beta = 1/T) = -\exp[i\theta_f]q_f(x, 0)$$



QCD with FTBC is  $Z_3$  invariant.

### PNJL Model with FTBC

$$D_{\nu} = \partial_{\nu} + iA_{\nu} = \partial_{\nu} + i\delta_{\nu,4}A_{4,a}\tilde{\lambda}_{a}/2$$
 PNJL Lagrangian quark part (Nambu-Jona-Lasinio type)

$$\begin{split} \mathcal{L} &= \sum_{f} \bar{q}_f (\gamma_\nu D_\nu - \mu_f \gamma_4 + m_f) q_f \\ &- G_{\mathrm{S}} \sum_{f} \sum_{a=0}^{8} [(\bar{q}_f \lambda_a q_f)^2 + (\bar{q}_f i \gamma_5 \lambda_a q_f)^2] \\ &+ G_{\mathrm{D}} \left[ \det_{ij} \bar{q}_i (1 + \gamma_5) q_j + \det_{ij} \bar{q}_i (1 - \gamma_5) q_j \right] \\ &+ \mathcal{U}(\varPhi[A], \bar{\varPhi}[A], T), \quad \underline{\qquad} \text{gluon potential} \end{split}$$

**FTBC** 

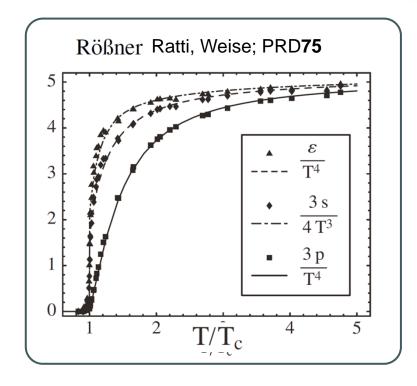
$$q_f(x, \beta = 1/T) = -\exp[i\theta_f]q_f(x, 0)$$

$$\Rightarrow \Phi = \frac{1}{3} \operatorname{Tr}_{c} e^{-iA_{4}/T}$$

## Polyakov potential

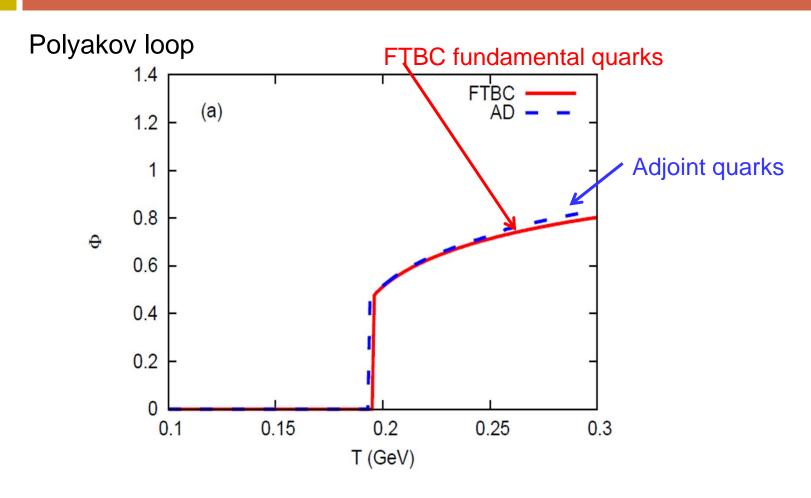
$$\mathcal{U} = T^4 \left[ -\frac{a(T)}{2} \Phi^* \Phi + b(T) \ln(1 - 6\Phi \Phi^* + 4(\Phi^3 + \Phi^{*3}) - 3(\Phi \Phi^*)^2) \right],$$

$$a(T) = a_0 + a_1 \left( \frac{T_0}{T} \right) + a_2 \left( \frac{T_0}{T} \right)^2, \quad b(T) = b_3 \left( \frac{T_0}{T} \right)^3.$$



It reproduces the lattice data in the pure gauge limit.

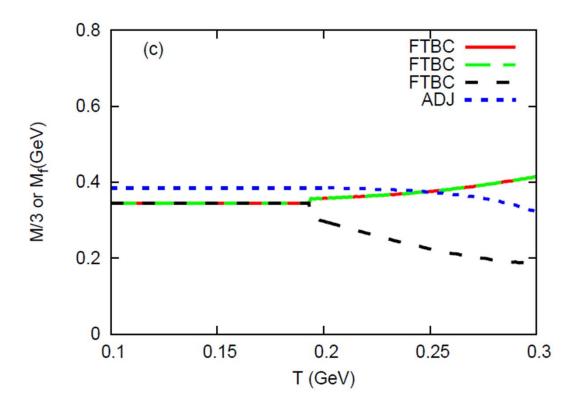
# PNJL with FTBC and adjoint quarks



FTBC and adjoint quarks preserve  $Z_3$  symmetry and hence the results are similar to each other for the Polyakov loop.

## PNJL with FTBC and adjoint quarks

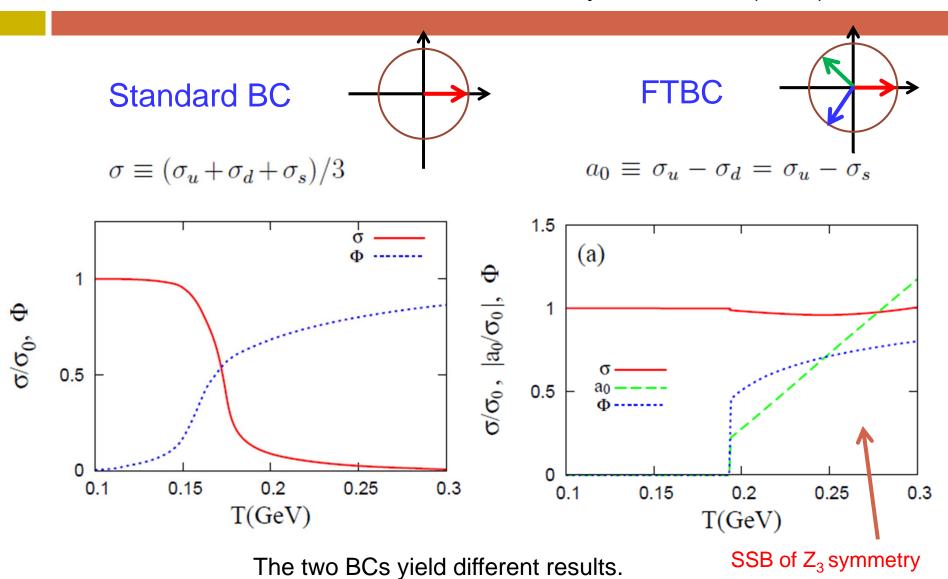
Dynamical quark mass (chiral condensate)



FTBC and adjoint quarks yield similar results also for chiral condensate at T smaller than the critical temperature.

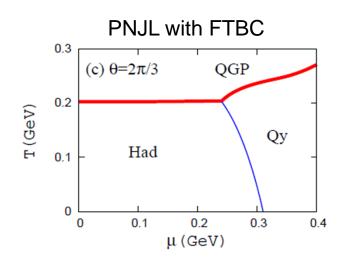
#### PNJL with the standard BC and FTBC

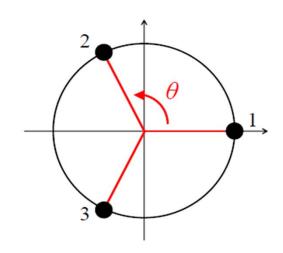
From Phys.Lett. B718 (2012) 130-135



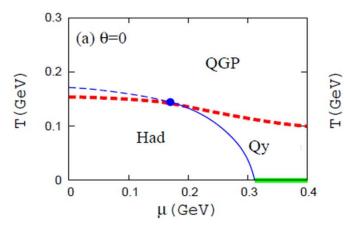
## Phase diagram

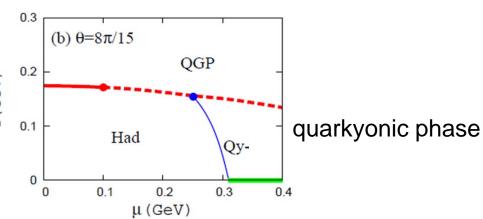
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PNJL with standard BC





# Summary

- 1. FTBC fundamental and adjoint quarks yield similar results, since the two fermions preserve  $Z_3$  symmetry.
- Fundamental and FTBC quarks yield rather different results.
- Therefore, the essential difference between fundamental and adjoint quarks comes from the presence or absence of Z<sub>3</sub> symmetry.

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